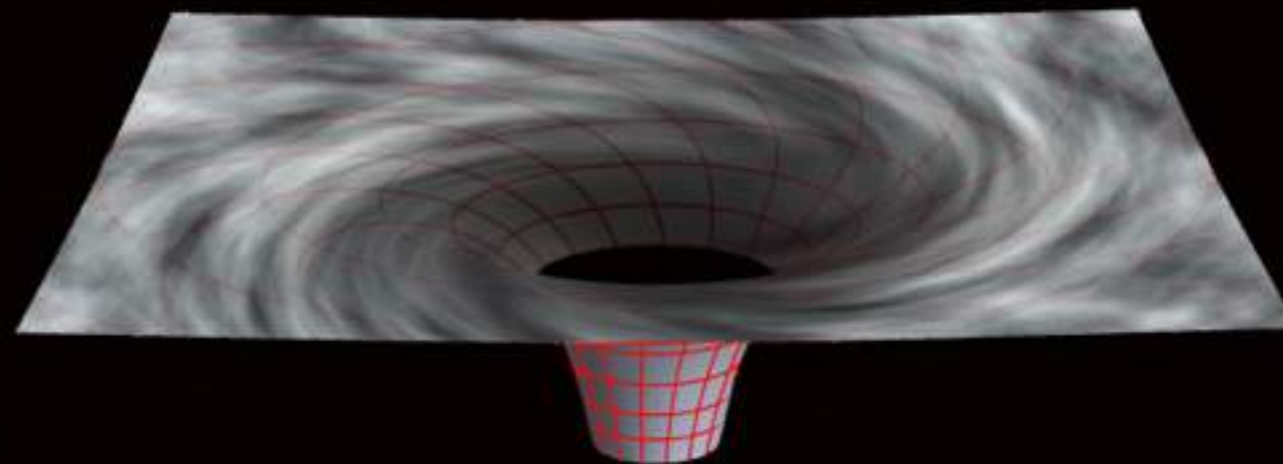


Rotating shells, non-spherical gravitational collapse and cosmic censorship

Jorge V. Rocha (Centra-IST, U.Lisboa)



- ❖ **JVR** Int. J. Mod. Phys. D24, 1542002 (2015) [1501.06724 [gr-qc]]
- ❖ **T. Delsate, JVR and R. Santarelli** Phys. Rev. D89, 121501(R) (2014) [1405.1433 [gr-qc]]
- ❖ **JVR, R. Santarelli, and T. Delsate** Phys. Rev. D89, 104006 (2014) [1402.4161 [gr-qc]]

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- ◆ There is overwhelming observational evidence that black holes (BHs) exist.

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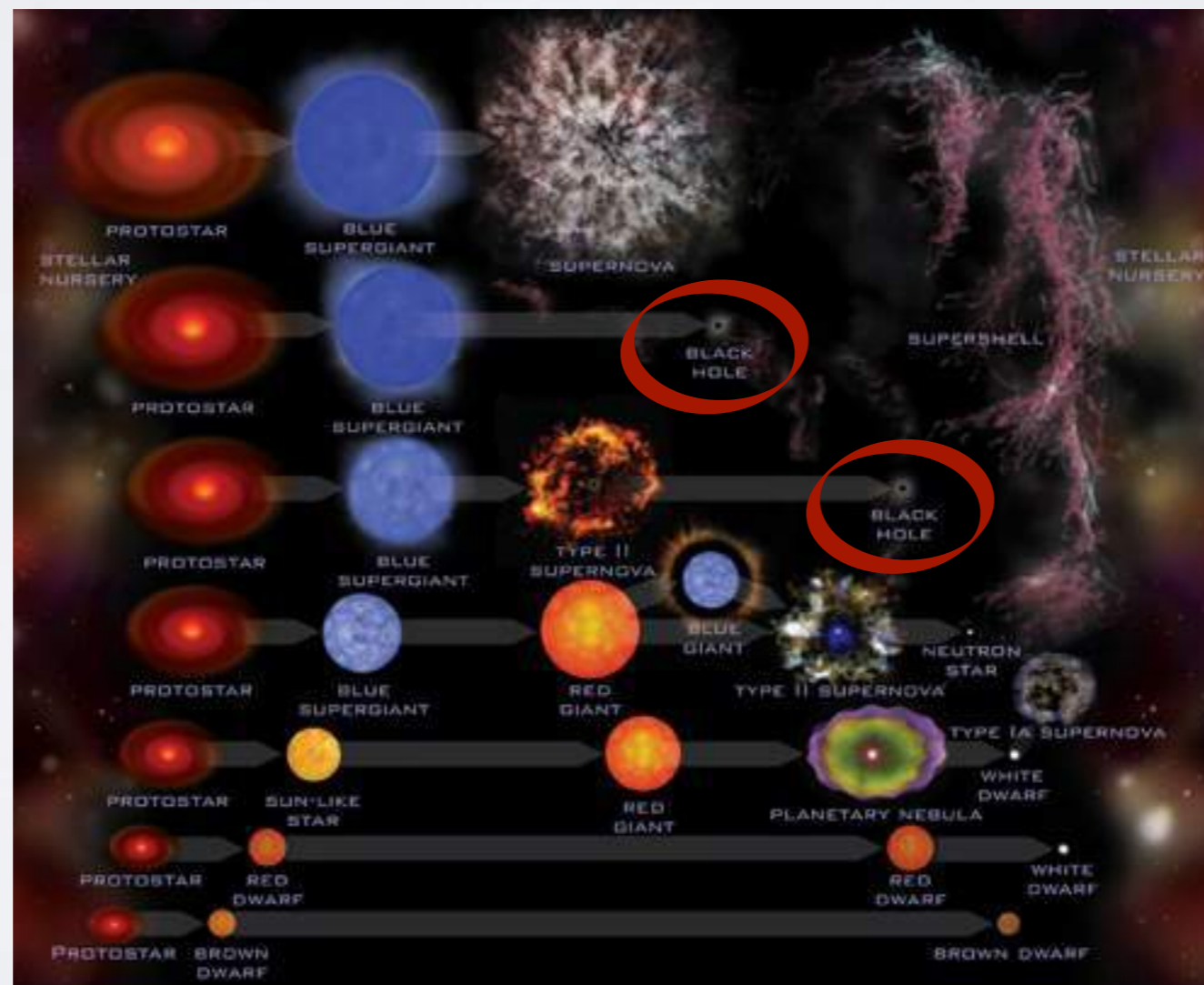
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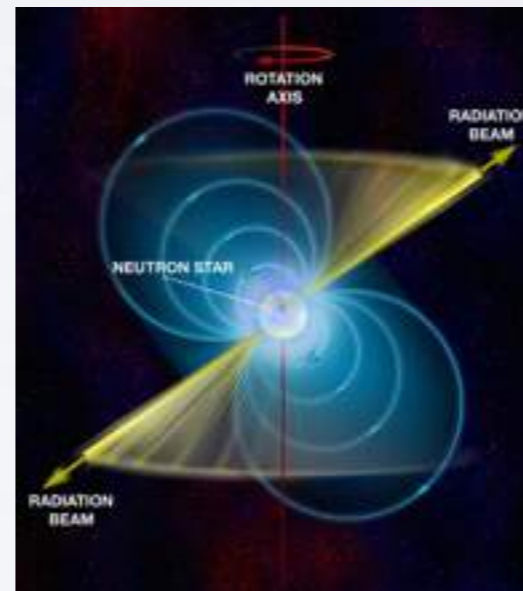
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- ◆ There is overwhelming observational evidence that black holes (BHs) exist.
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- ◆ BHs are theoretically predicted as the endpoint of gravitational collapse of sufficiently massive stars.
- ◆ The vast majority of celestial objects are rotating. Black holes are no exception.



ESO / J. Pérez



ESO

Introduction: Gravitational collapse with rotation

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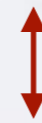
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1. realistic collapses should include rotation;
2. known 'violations' of the cosmic censorship conjecture (CCC) occur in non-rotating — thus non-generic — settings;
3. rotation introduces instabilities (e.g. superradiance).

Introduction: Approaching the problem

- ◆ Advantage of non-rotating setups is their large amount of symmetry. Spherical symmetry reduces problem to 1+1 dims.

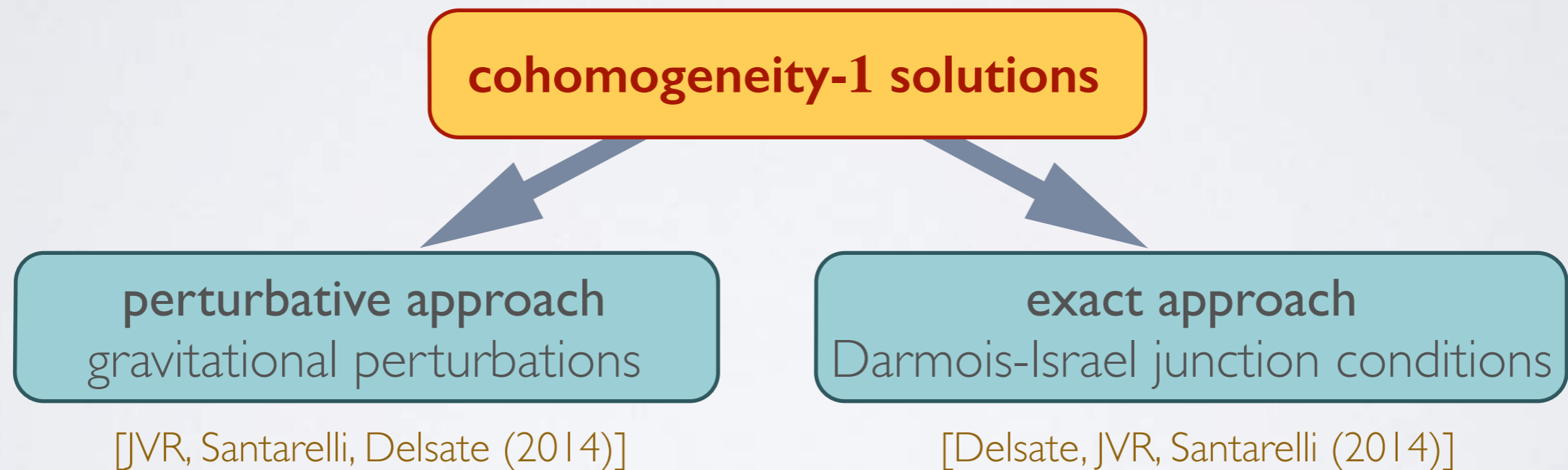
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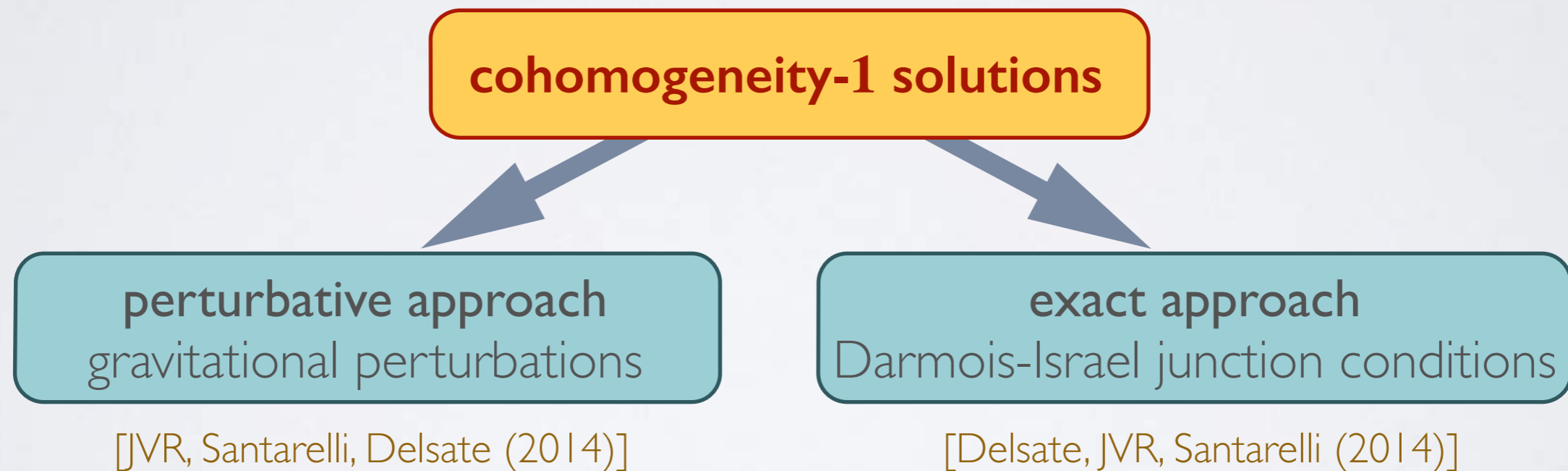
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- ♦ The price to pay for the convenience provided by cohomogeneity-1 spacetimes is the restriction to higher (odd) dimensions, $D=2N+3$ with $N=1, 2, 3, \dots$

Background: Cohomogeneity-1 black holes

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- ◆ Constant t and r sections are squashed $(2N+1)$ -spheres.
- ◆ S^{2N+1} can be written as a S^1 bundle over CP^N .

Background: Cohomogeneity-1 black holes

- ◆ The metric for these cohomogeneity-1 BHs is

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu = -f(r)^2 dt^2 + g(r)^2 dr^2 + r^2 \hat{g}_{ab} dx^a dx^b + h(r)^2 [d\psi + A_a dx^a - \Omega(r) dt]^2$$

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where

$$g(r)^2 = \left(1 + \frac{r^2}{\ell^2} - \frac{2M\Xi}{r^{2N}} + \frac{2Ma^2}{r^{2N+2}} \right)^{-1}, \quad f(r) = \frac{r}{g(r)h(r)},$$
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\hat{g}_{ab} denotes the Fubini-Study metric on CP^N and $A_a dx^a$ is its Kahler potential.

For $N=1$:

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- ◆ n.b. These solutions accommodate a non-vanishing cosmological constant:

$$R_{\mu\nu} = -(D-1)\ell^{-2}g_{\mu\nu}$$

Background: Thin shells in cohomogeneity-1 spacetimes

- ◆ The cohomogeneity-1 property makes an exact (thin-shell) calculation possible, 'gluing' an interior to an exterior geometry.

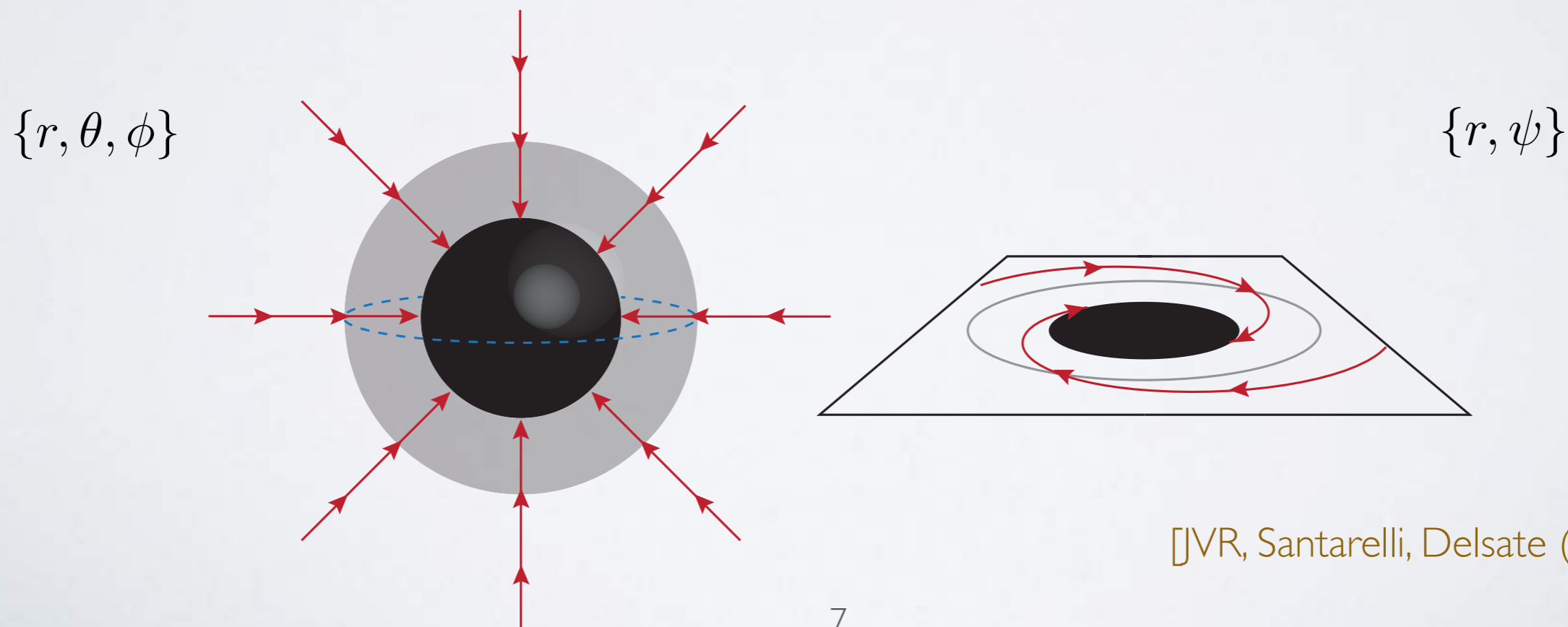
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- ◆ n.b. The dynamics on the $CP^1 \cong S^2$ and on the S^1 **separate**. All traces of the rotation show up in the $\{r, \psi\}$ plane.



[JVR, Santarelli, Delsate (2014)]

Rotating thin shells: **junction conditions**

- ◆ Use junction conditions along a timelike hypersurface, $t = \mathcal{T}(\tau)$, $r = \mathcal{R}(\tau)$:

$$\mathfrak{g}_{ij}^{(+)} = \mathfrak{g}_{ij}^{(-)} \equiv \mathfrak{g}_{ij} ,$$
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energy density \longleftarrow ρ pressure \longleftarrow P intrinsic momentum / heat flow \longleftarrow 2φ pressure anisotropy \longleftarrow ΔP

Rotating thin shells: Equation of state and shell equation of motion

- ◆ The stress-energy tensor components are dictated by the metric components:

$$\rho = -\frac{(\beta_+ - \beta_-)(\mathcal{R}^2 h)'}{8\pi\mathcal{R}^3}$$

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- ◆ These equations can be integrated, yielding the **shell's equation of motion**:

$$\dot{\mathcal{R}}^2 + V_{\text{eff}}(\mathcal{R}) = 0$$

Rotating thin shells: Effective potential

- ♦ For generic values of N , and a linear equation of state:

$$\dot{\mathcal{R}}^2 + V_{\text{eff}}(\mathcal{R}) = 0 \quad V_{\text{eff}}(\mathcal{R}) = 1 + \frac{\mathcal{R}^2}{\ell^2} + \frac{2Ma^2}{\ell^2 \mathcal{R}^{2N}} + \frac{2Ma^2}{\mathcal{R}^{2N+2}} - \frac{M_+ + M_-}{\mathcal{R}^{2N}} - \left(\frac{M_+ - M_-}{m_0} \right)^2 \left(\frac{\mathcal{R}^{2N}}{m_0} \right)^{\frac{2N+1}{N}w} \left(1 + \frac{2Ma^2}{\mathcal{R}^{2N+2}} \right)^{w-1} - \frac{1}{4} \left(\frac{m_0}{\mathcal{R}^{2N}} \right)^{2+\frac{2N+1}{N}w} \left(1 + \frac{2Ma^2}{\mathcal{R}^{2N+2}} \right)^{1-w} .$$

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$$\dot{\mathcal{R}}^2 + V_{\text{eff}}(\mathcal{R}) = 0 \quad V_{\text{eff}}(\mathcal{R}) = 1 + \frac{\mathcal{R}^2}{\ell^2} + \frac{2Ma^2}{\ell^2 \mathcal{R}^{2N}} + \frac{2Ma^2}{\mathcal{R}^{2N+2}} - \frac{M_+ + M_-}{\mathcal{R}^{2N}} \\ - \left(\frac{M_+ - M_-}{m_0} \right)^2 \left(\frac{\mathcal{R}^{2N}}{m_0} \right)^{\frac{2N+1}{N}w} \left(1 + \frac{2Ma^2}{\mathcal{R}^{2N+2}} \right)^{w-1} \\ - \frac{1}{4} \left(\frac{m_0}{\mathcal{R}^{2N}} \right)^{2+\frac{2N+1}{N}w} \left(1 + \frac{2Ma^2}{\mathcal{R}^{2N+2}} \right)^{1-w} .$$

- ◆ For $N=1$ and large values of \mathcal{R} :

$$V_{\text{eff}} \approx 1 + \frac{\mathcal{R}^2}{\ell^2} - \left(\frac{\Delta M}{m_0} \right)^2 \left(\frac{\mathcal{R}^2}{m_0} \right)^{3w} - \frac{1}{4} \left(\frac{m_0}{\mathcal{R}^2} \right)^{2+3w}$$

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$$V_{\text{eff}} \approx \frac{2Ma^2}{\mathcal{R}^4} - \frac{M_+ + M_-}{\mathcal{R}^2} - \frac{1}{4} \left(\frac{2Ma^2}{m_0^2} \right)^{1-w} \left(\frac{m_0}{\mathcal{R}^2} \right)^{4+w} - \left(\frac{2Ma^2}{m_0^2} \right)^{w-1} \left(\frac{\Delta M}{m_0} \right)^2 \left(\frac{\mathcal{R}^2}{m_0} \right)^{2+w}$$

Rotating thin shells: Stationary shell around a BH in AdS

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Rotating thin shells: Stationary shell around a BH in AdS

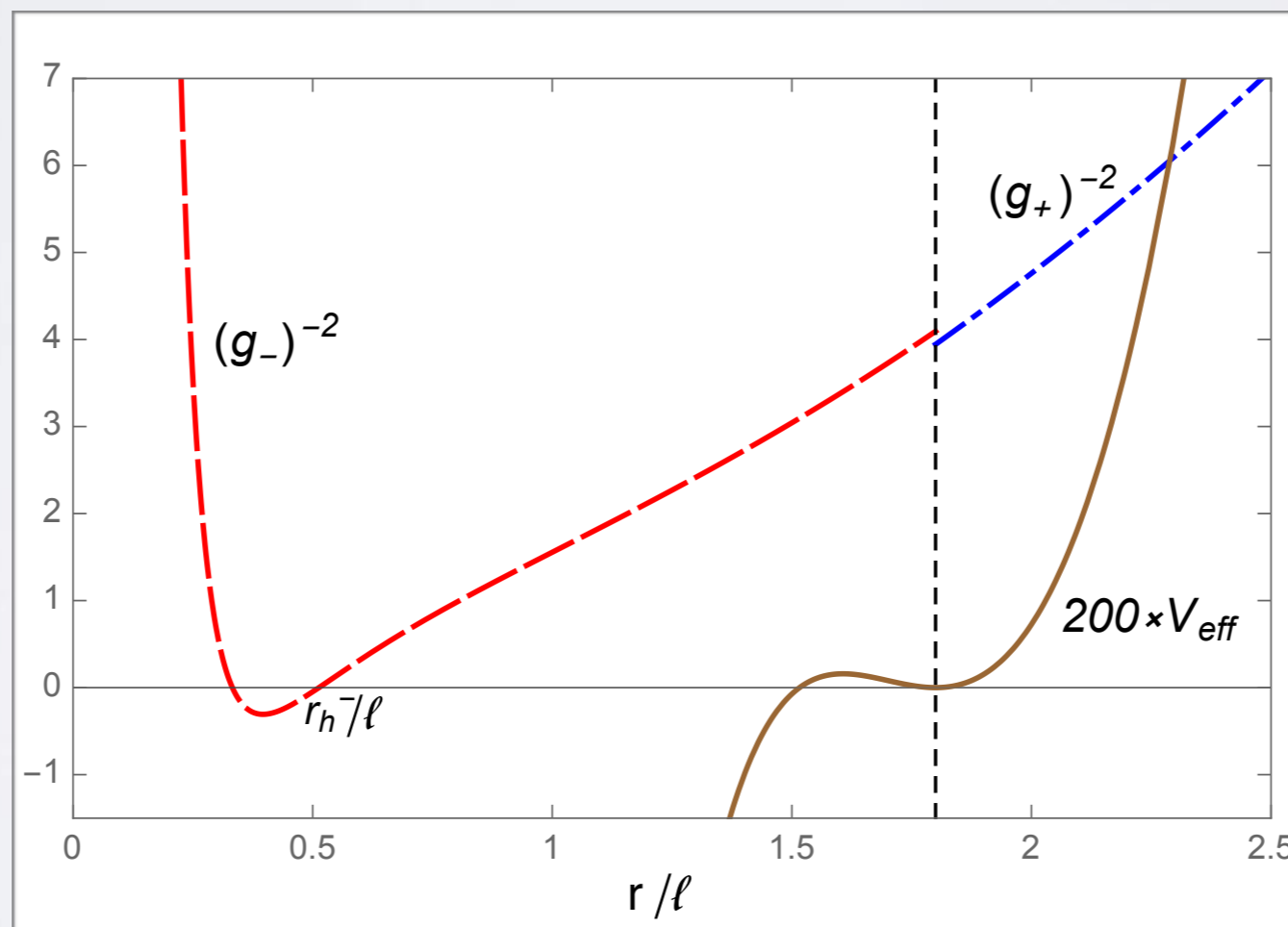
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$$m_0/\ell^2 = 0.324, \quad w = 0.285, \quad Ma^2/\ell^4 = 0.02, \quad \mathcal{R}_*/\ell = 1.8$$




[Delsate, JVR, Santarelli (2014)]


Rotating thin shells: Full collapse in asymptotically flat spacetime

◆ Take asymptotically flat limit, $\ell \rightarrow \infty$.

◆ Collapse starting from rest at infinity imposes:


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i.e., matter on the shell has EoS of dust


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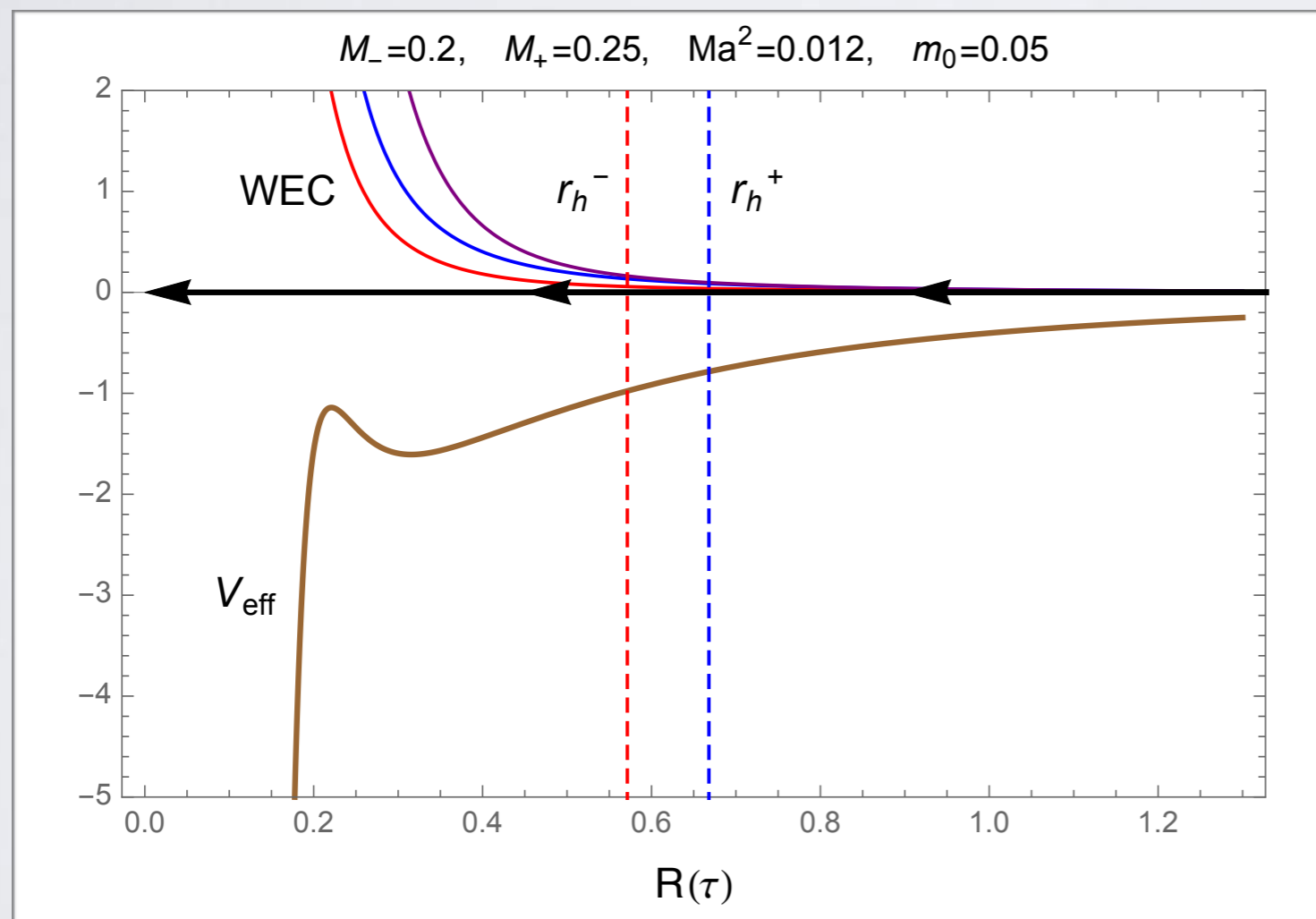
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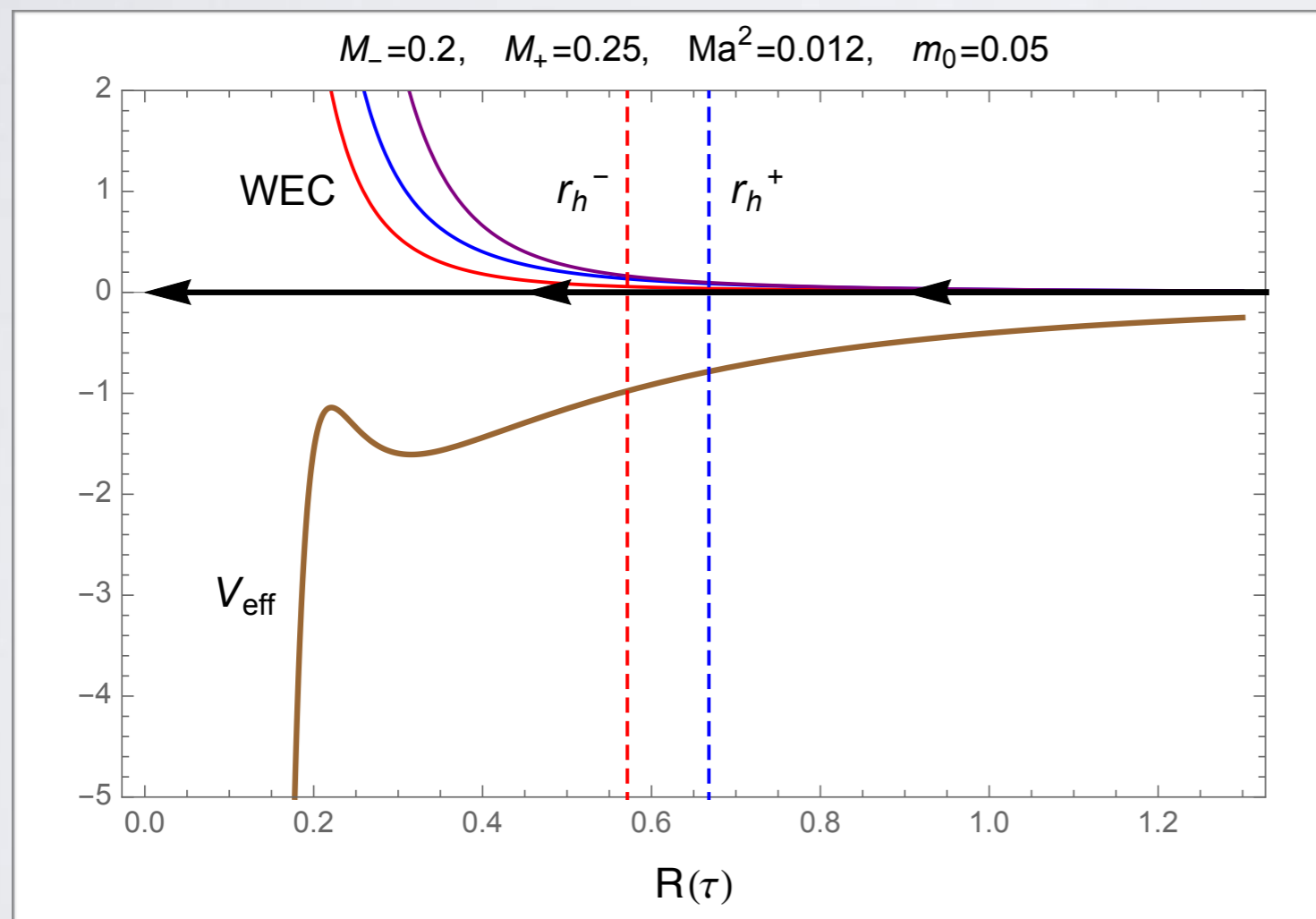
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Weak energy conditions (WEC) are satisfied

Rotating thin shells: Cosmic censorship

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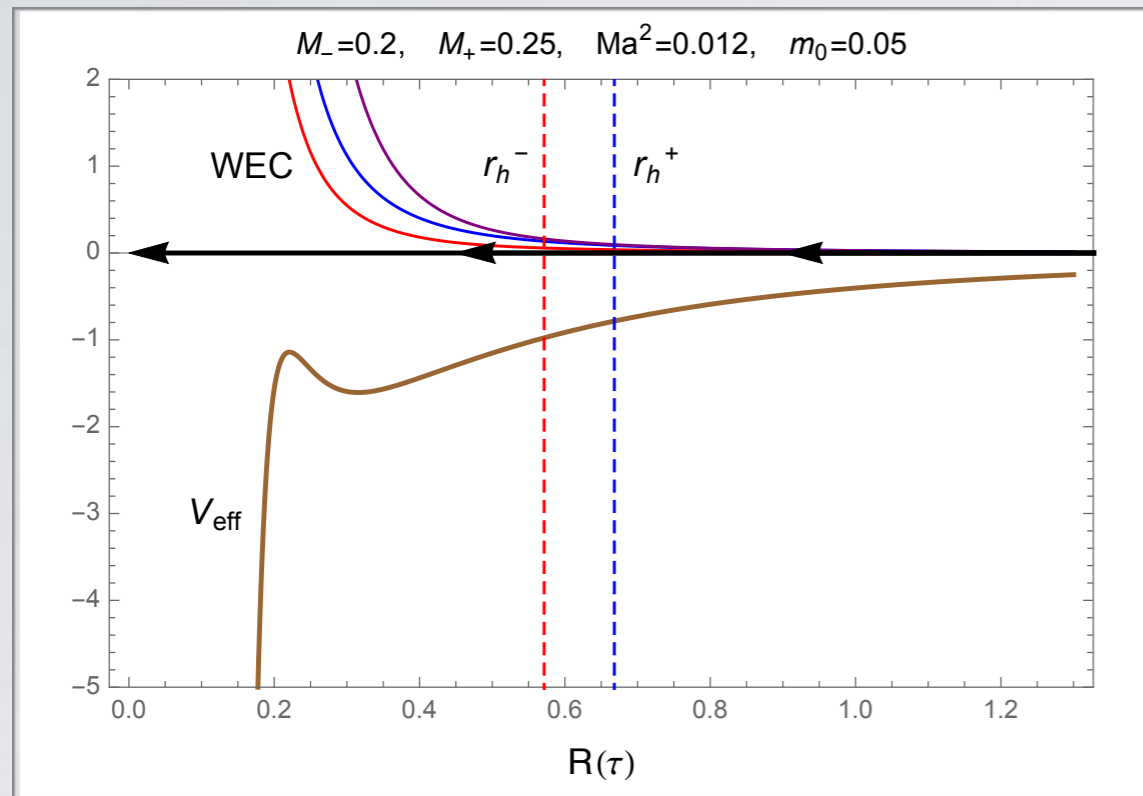
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CCC is preserved

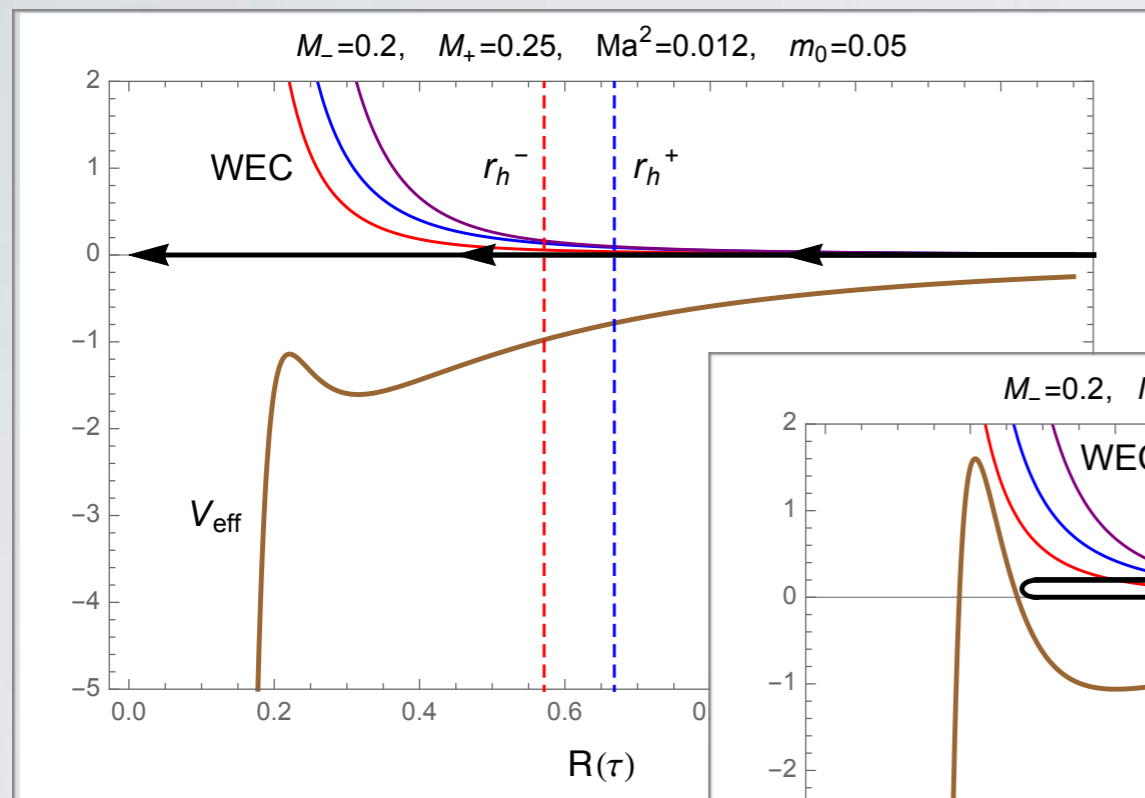
Rotating thin shells: Diverse scenarios



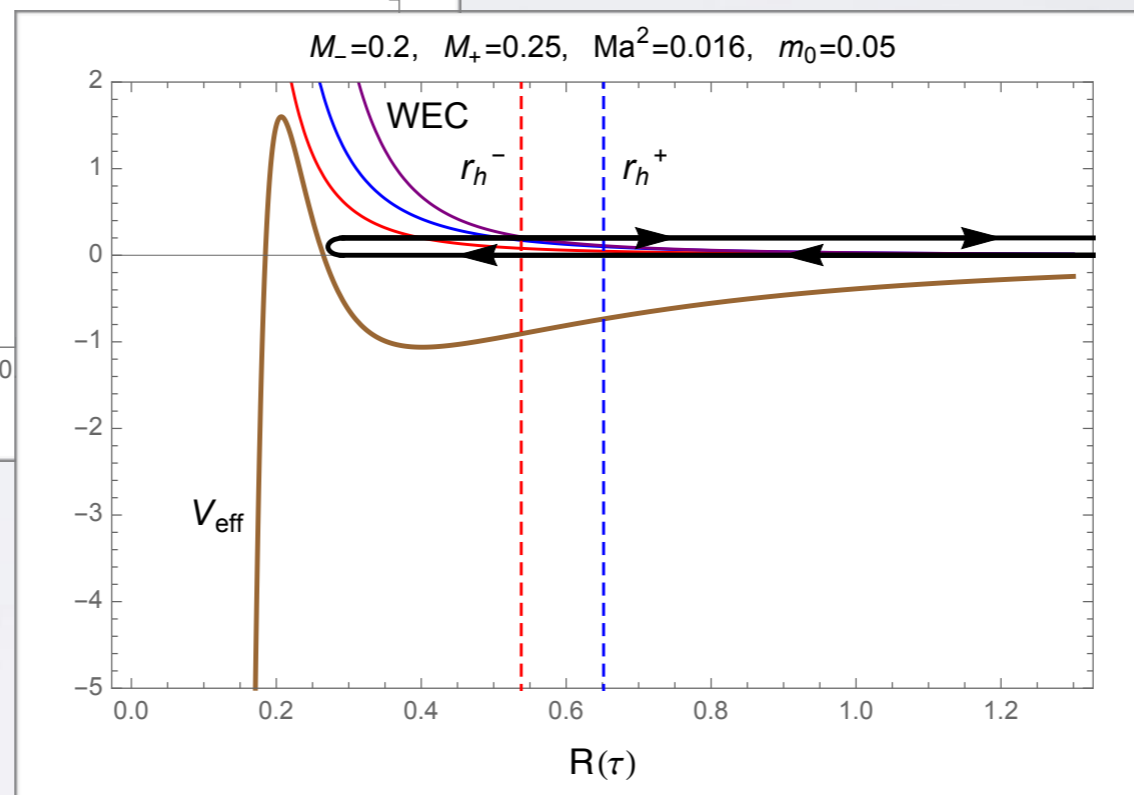
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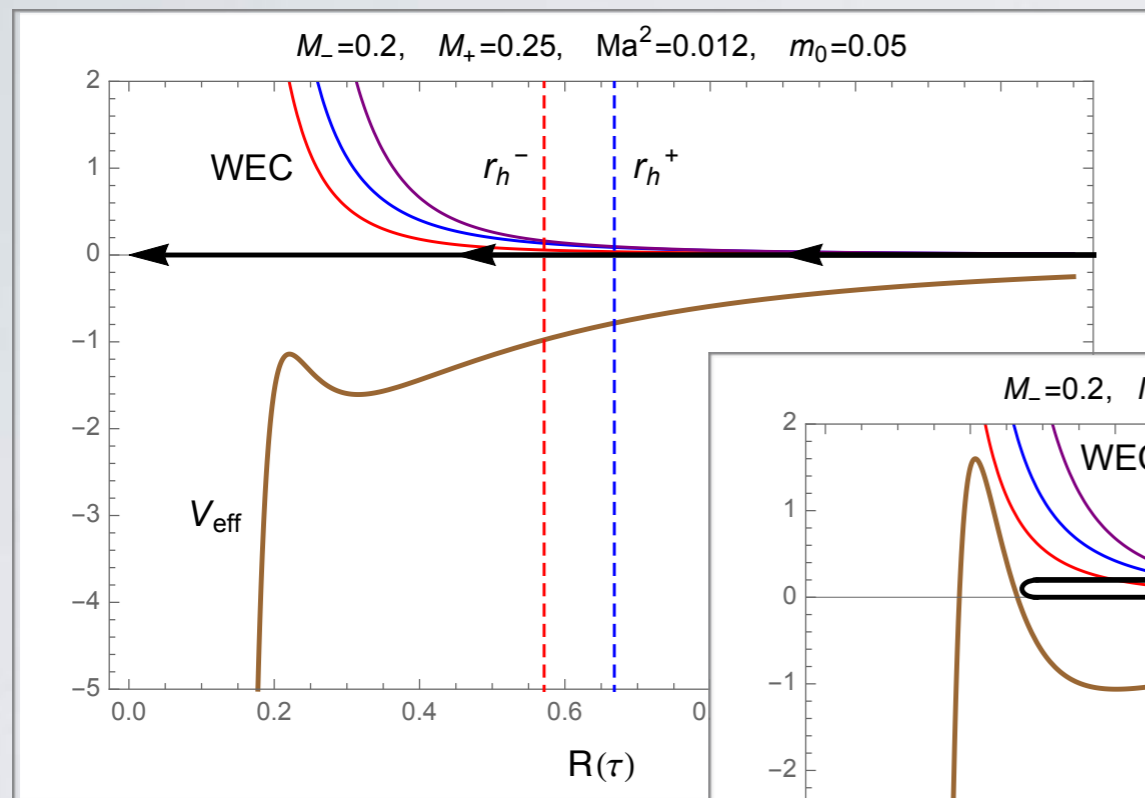
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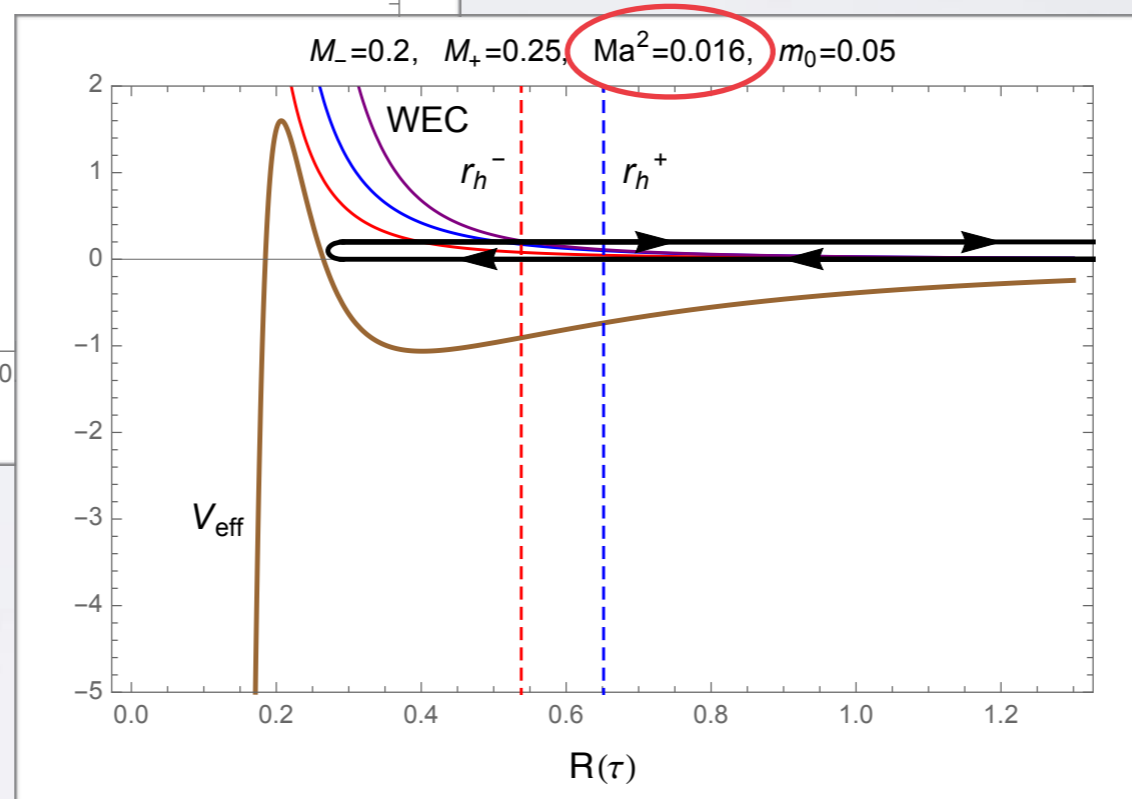
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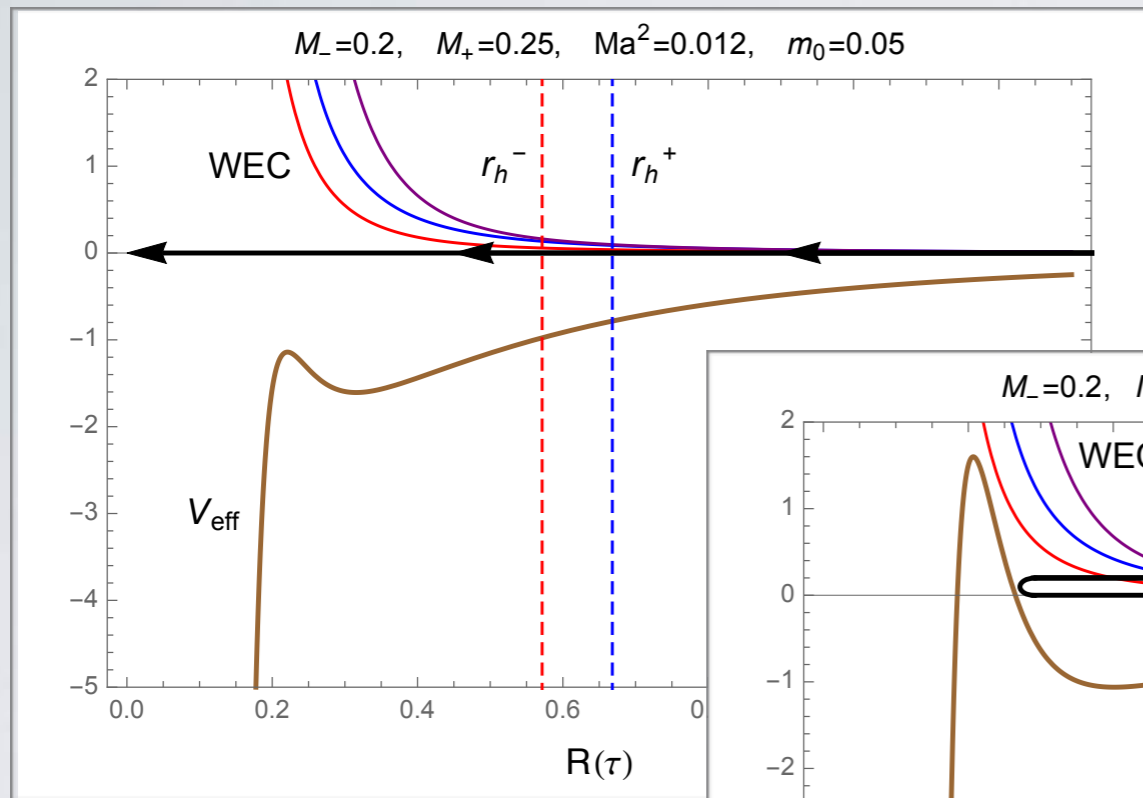
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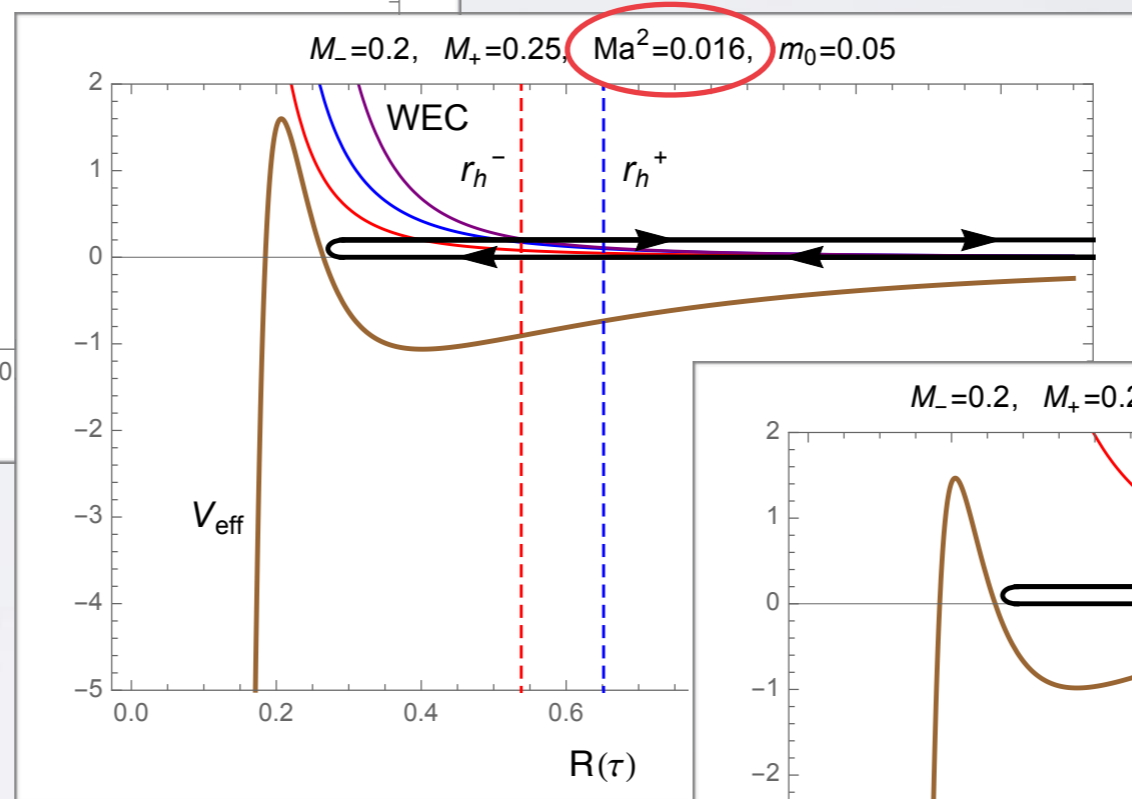
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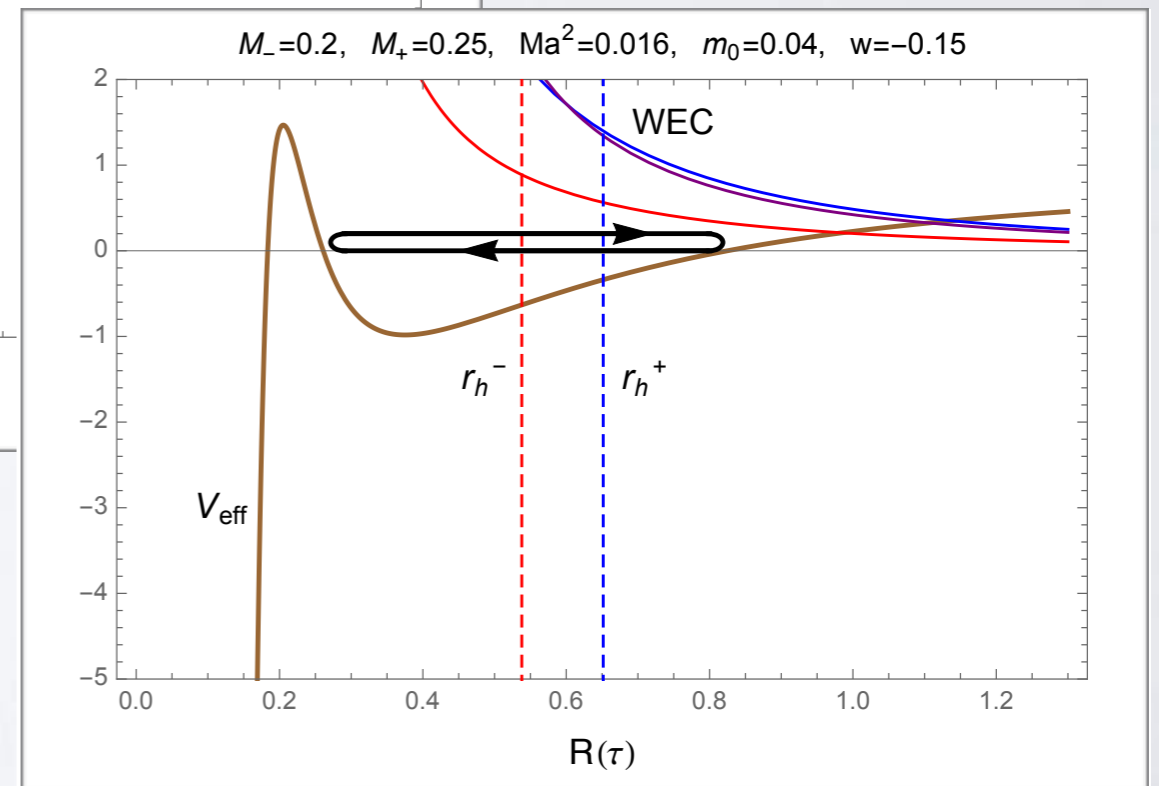


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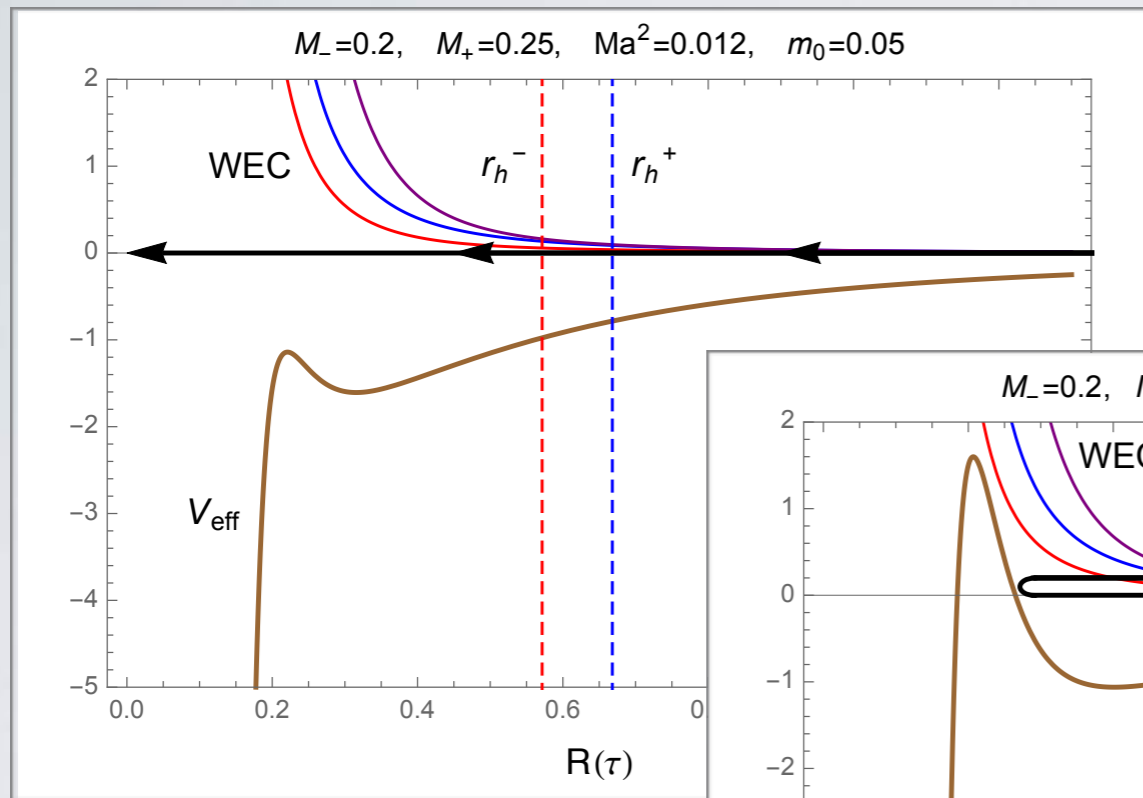


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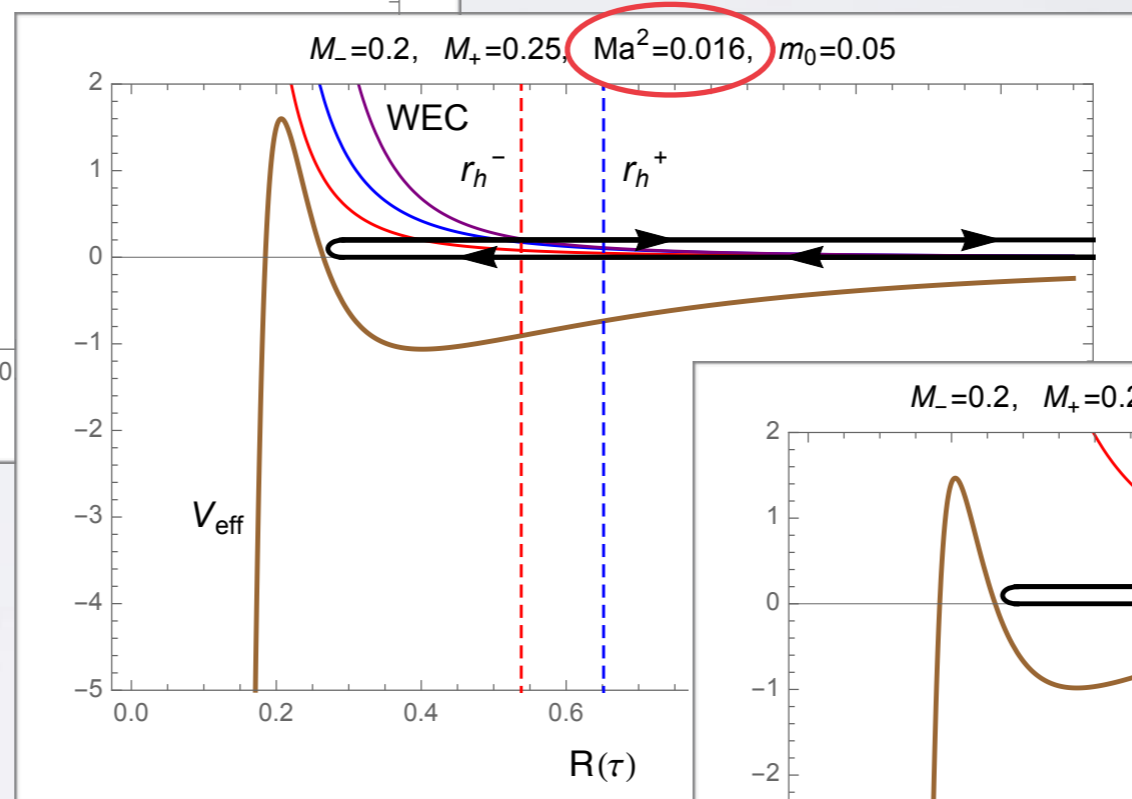
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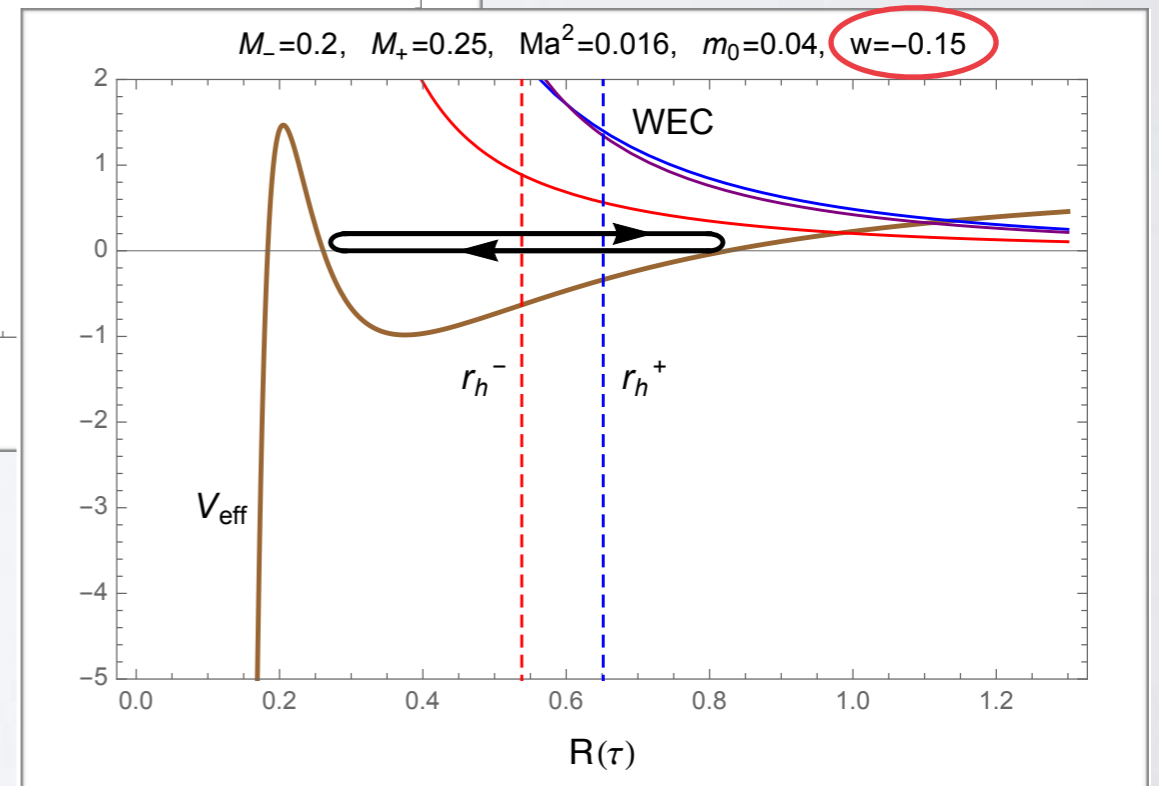


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Rotating thin shells: Scanning parameter space

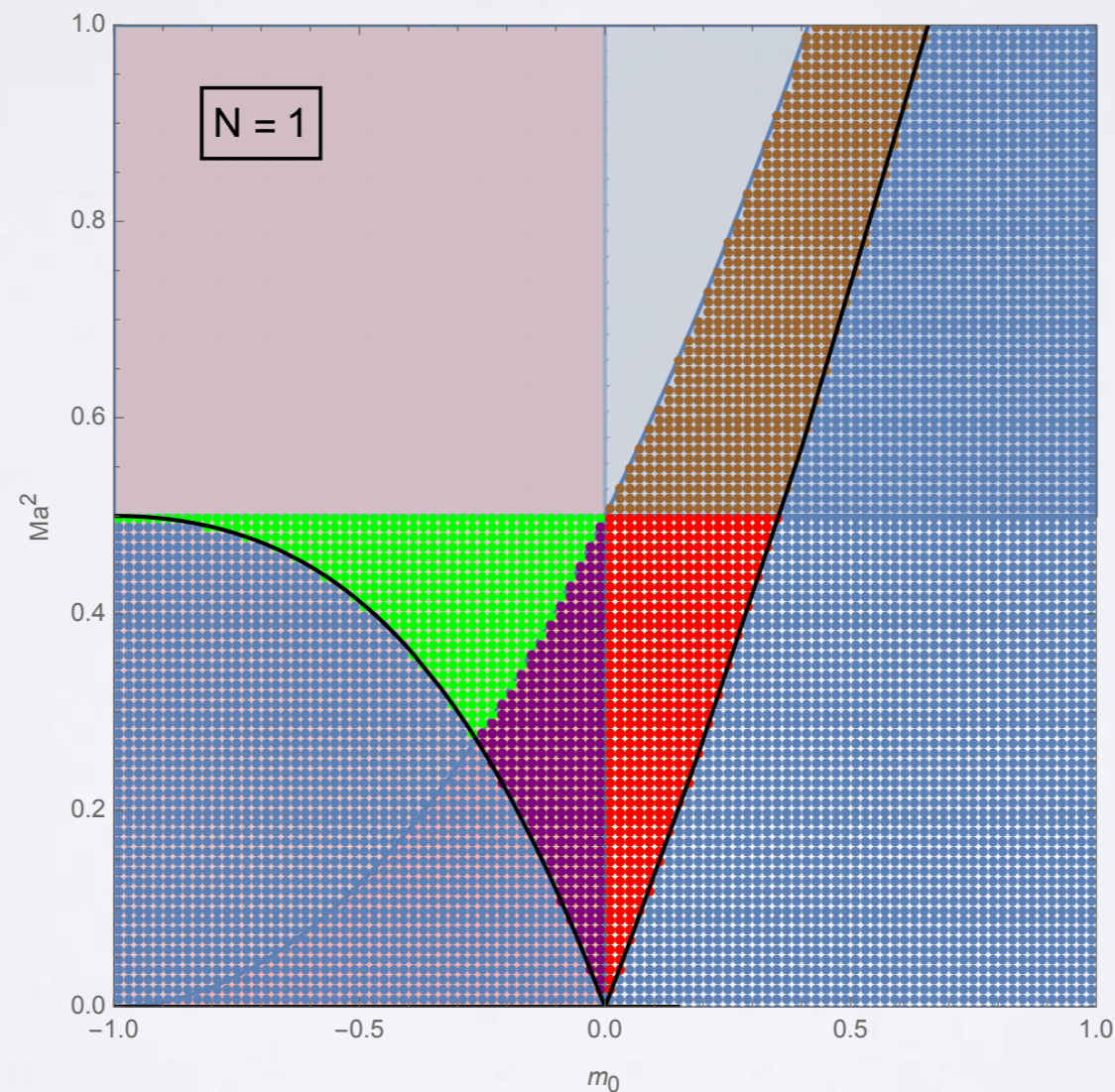
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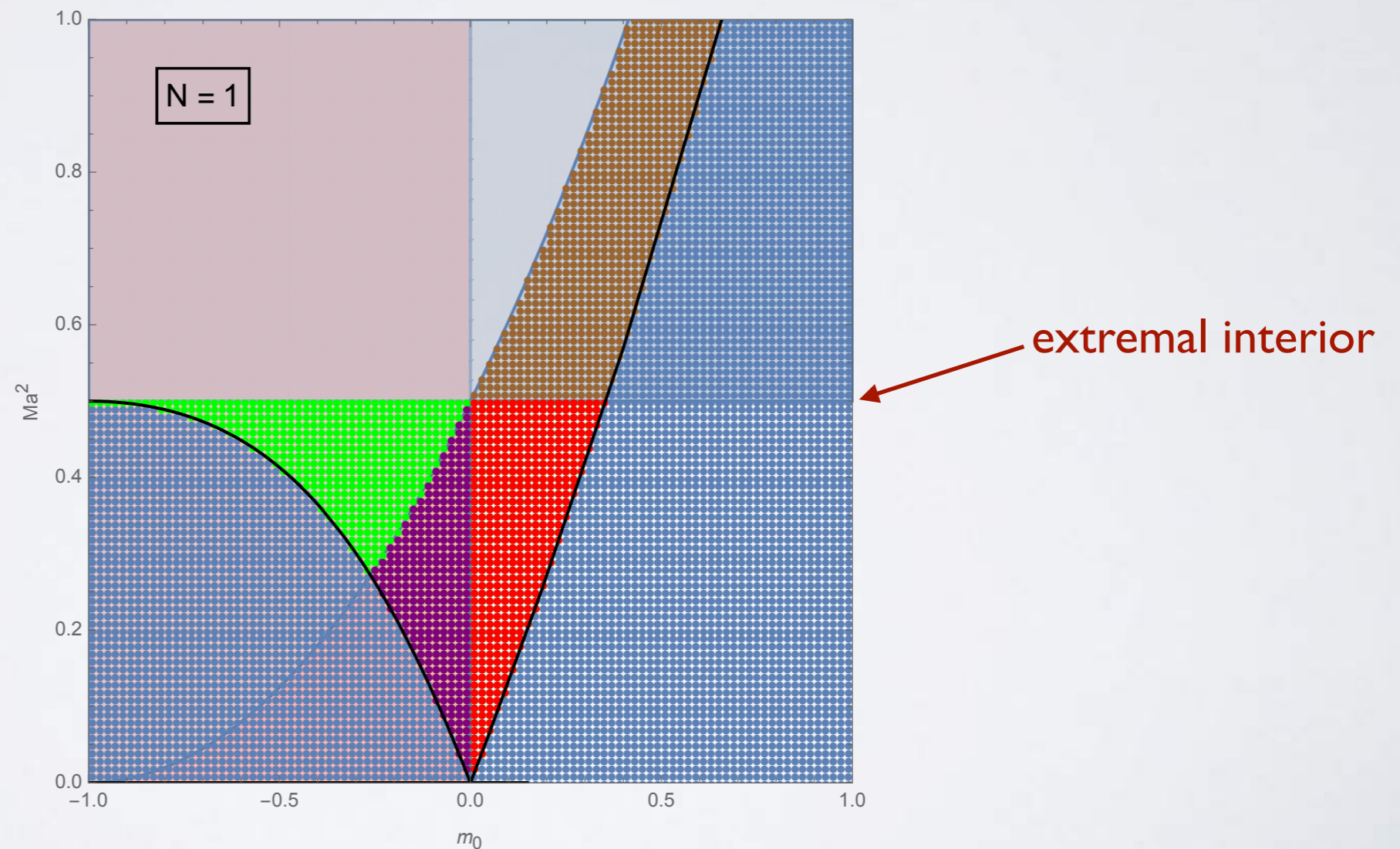
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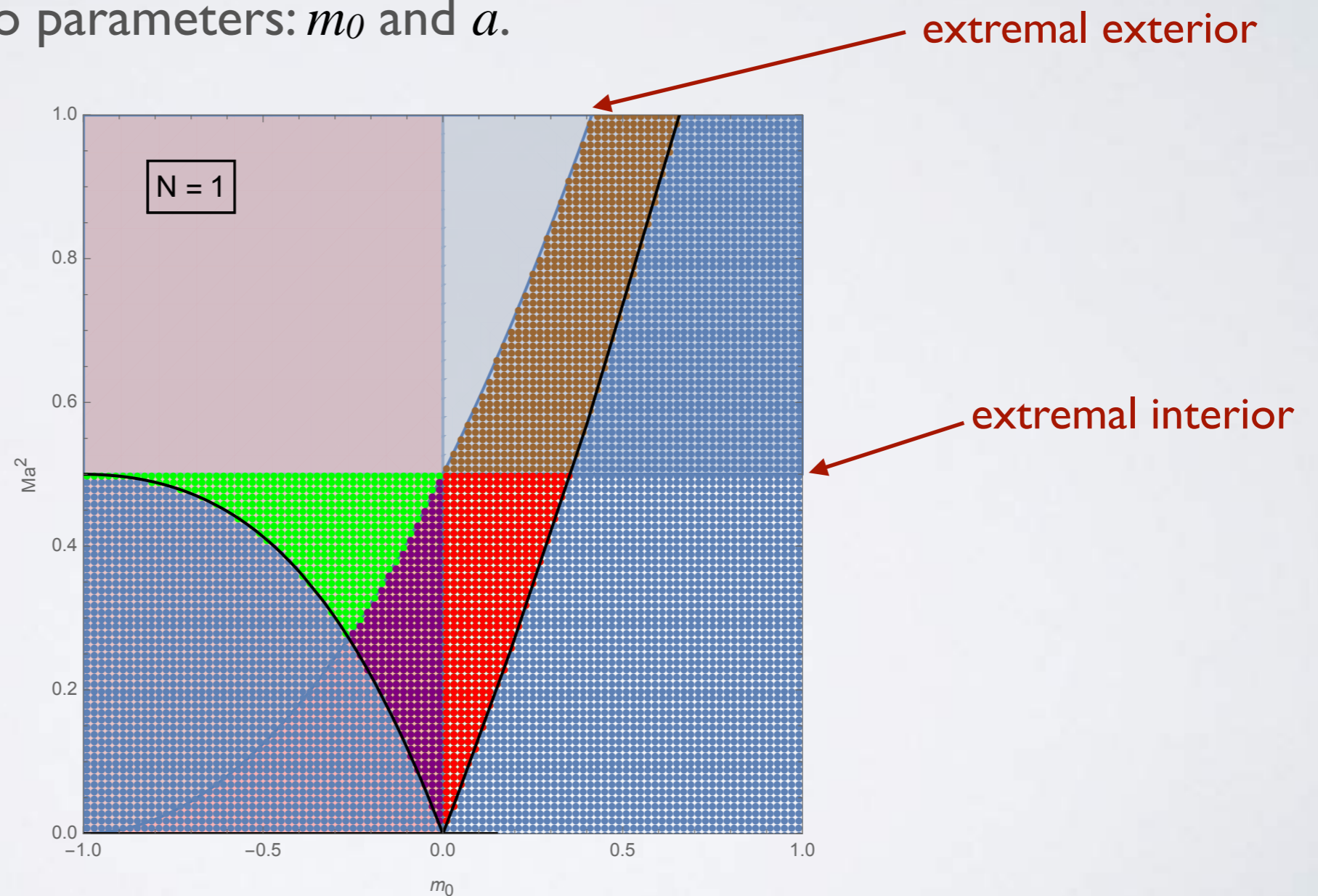
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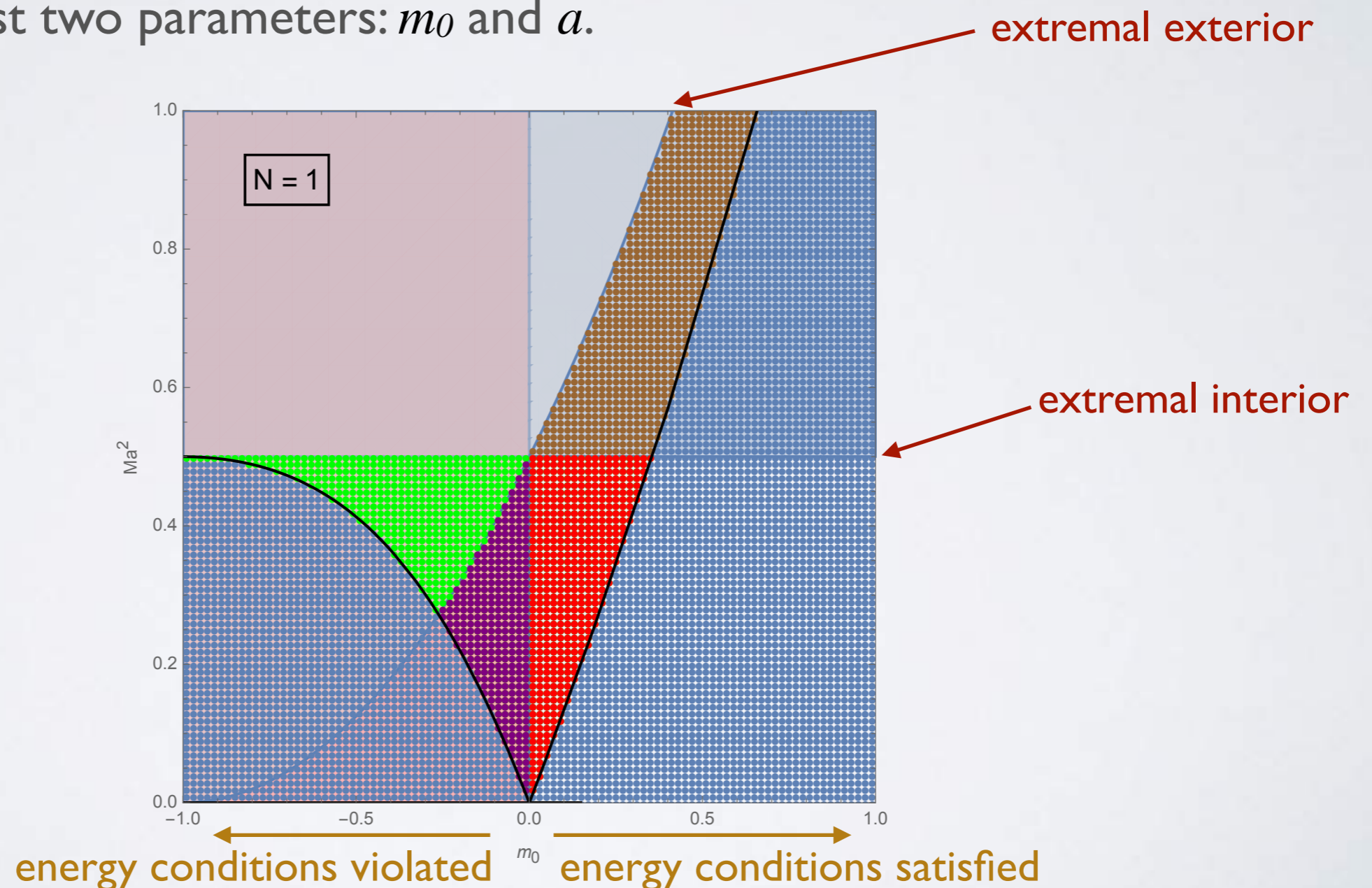
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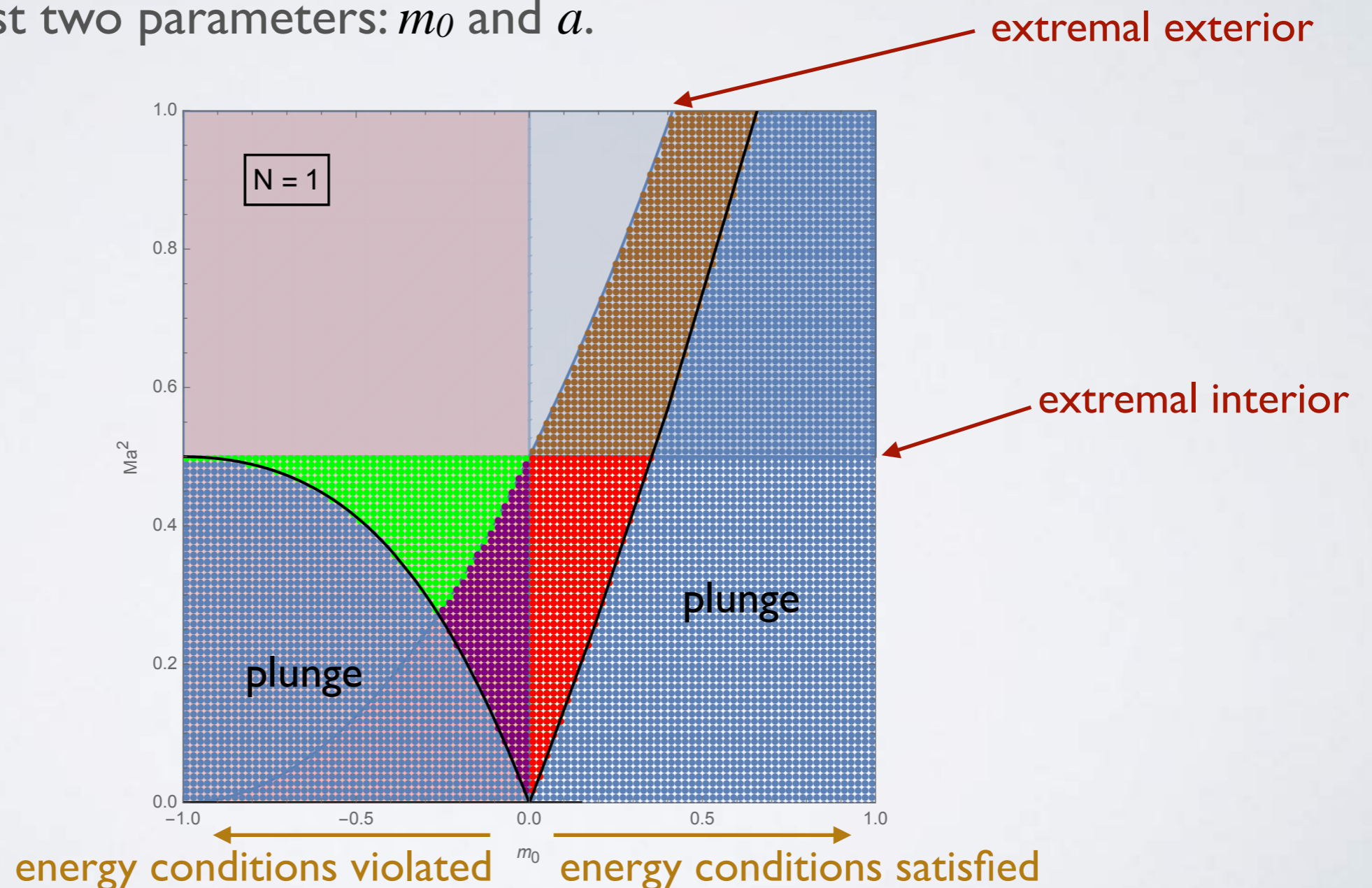
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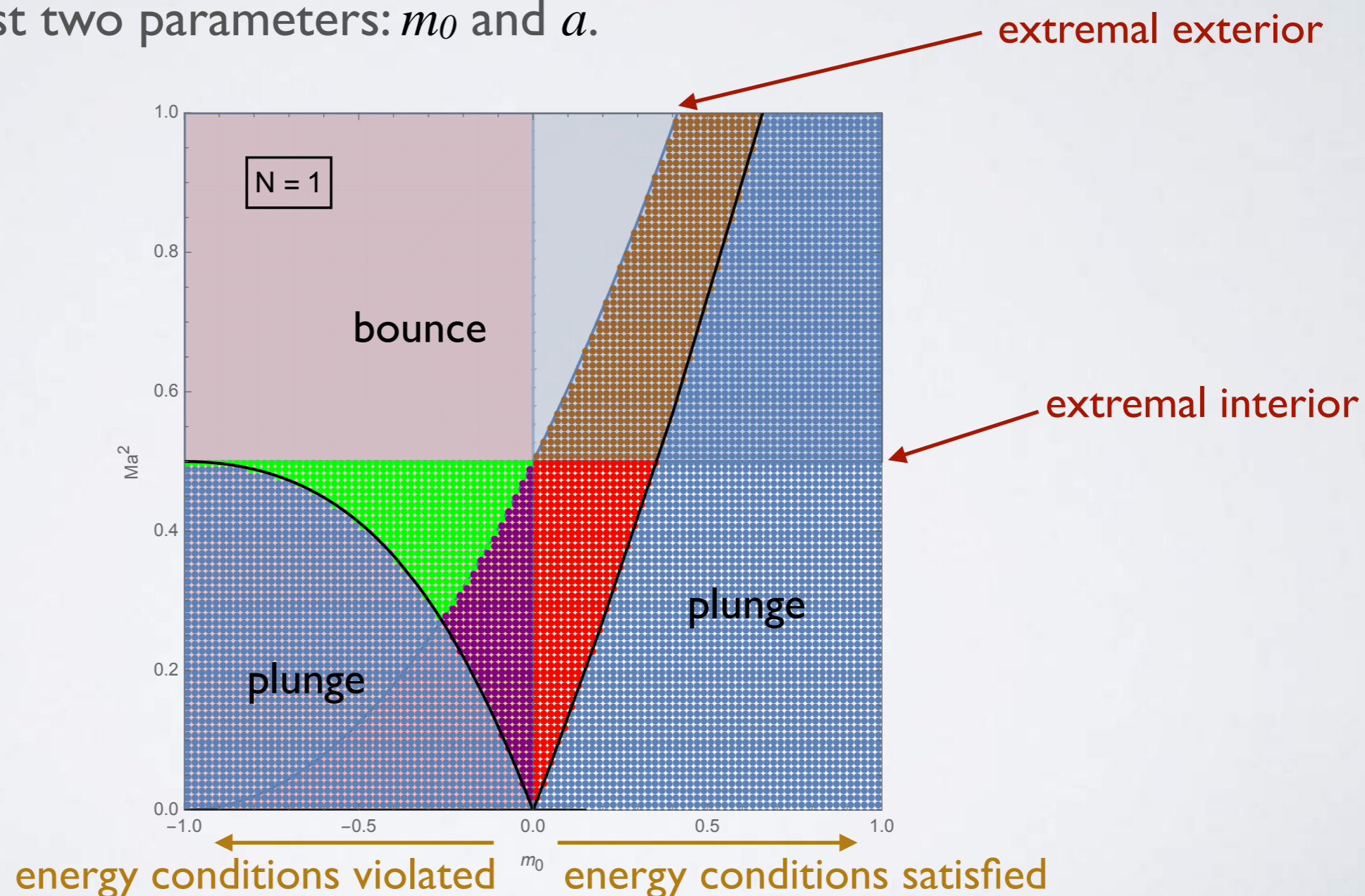
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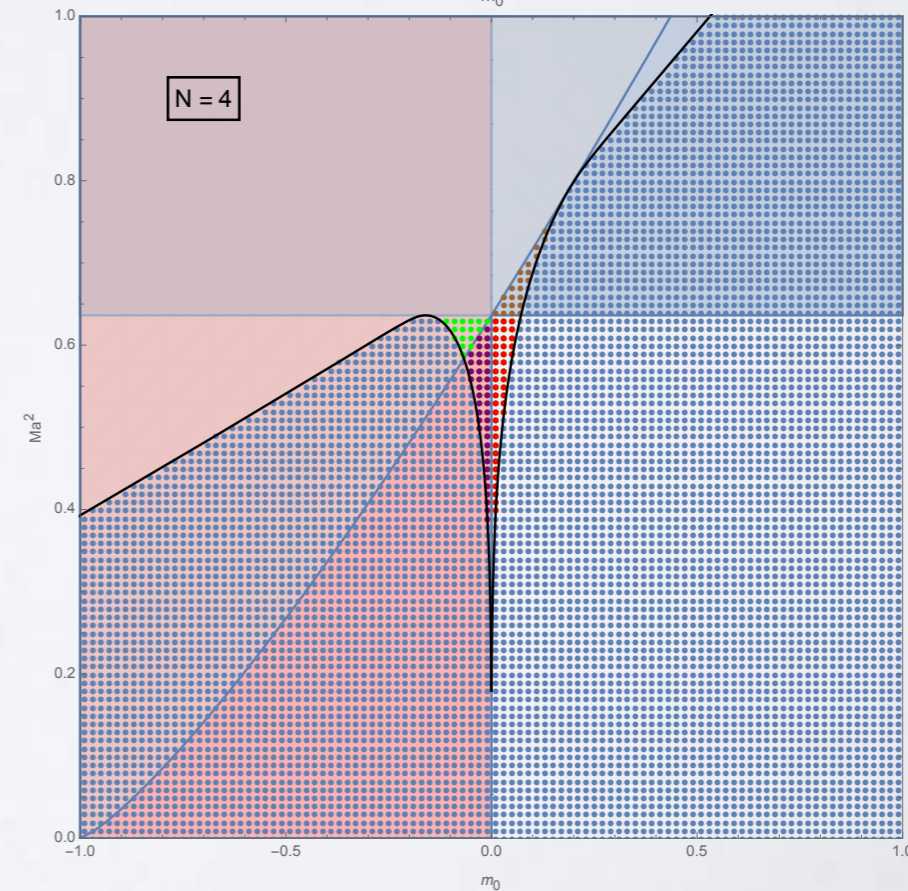
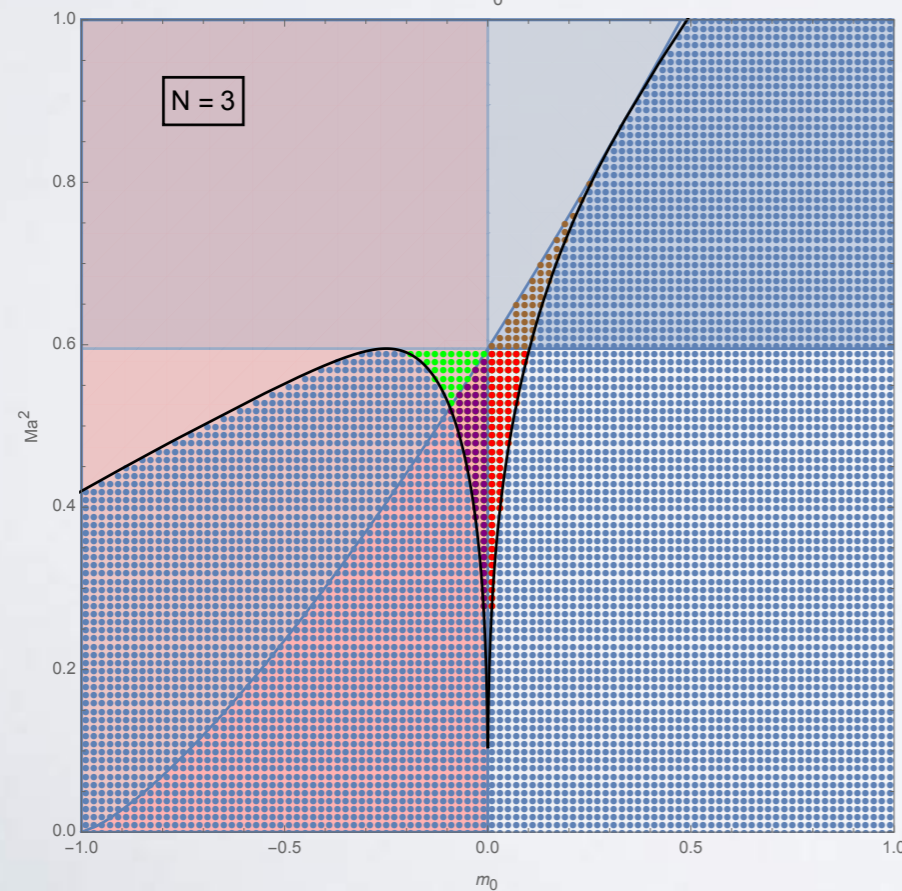
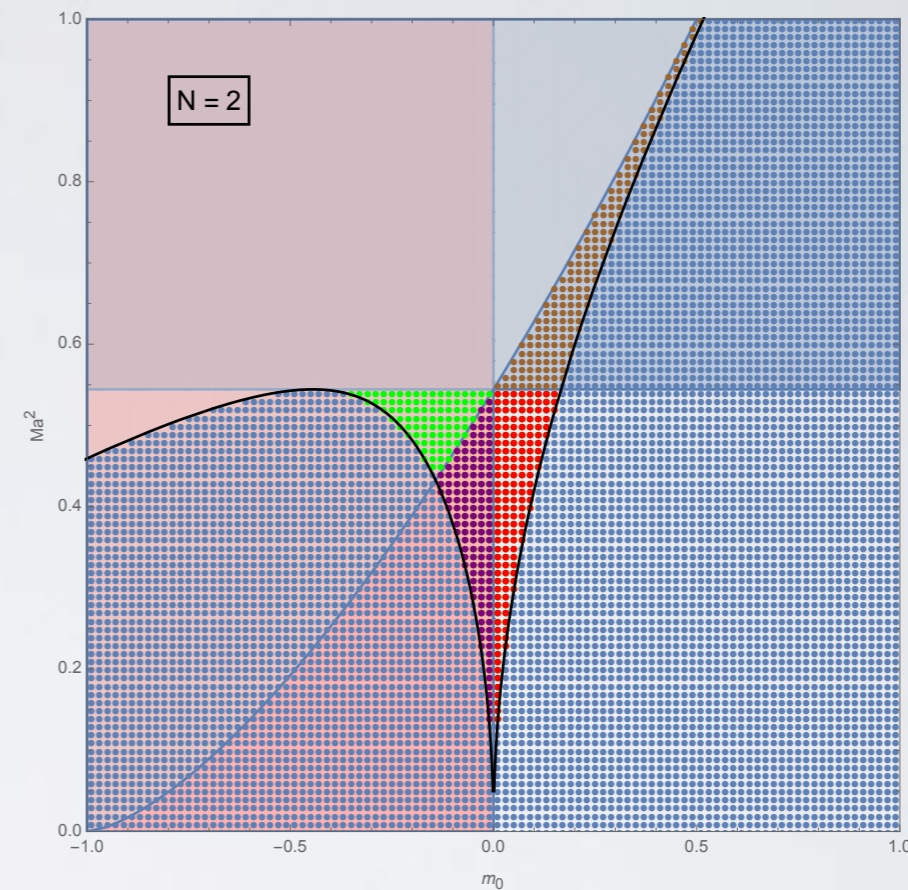
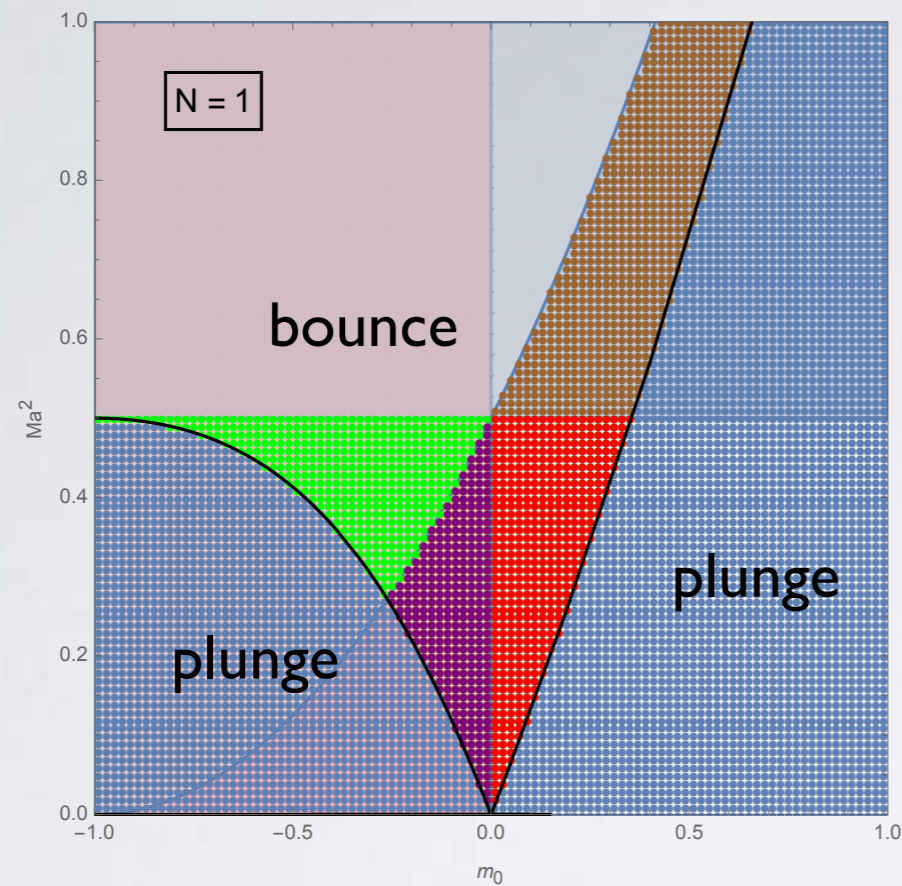


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